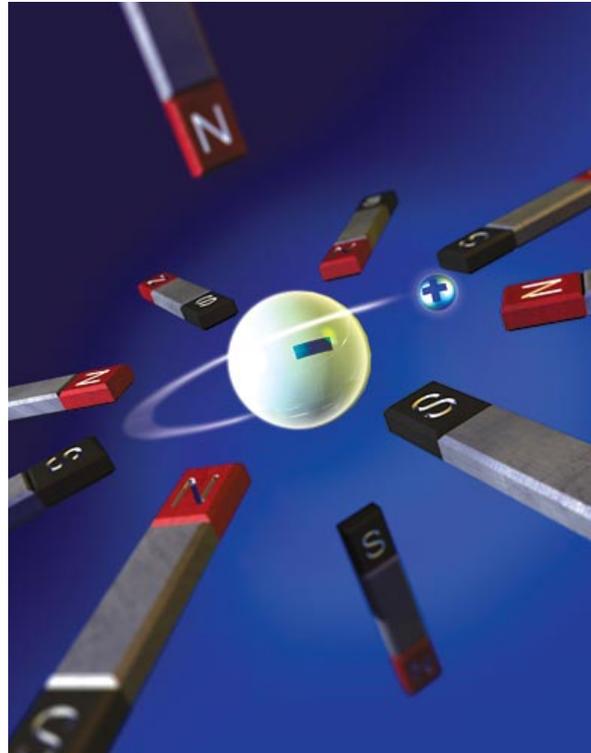


Properties of bound electrons and muons



PSAS'2016
Jerusalem, May 23, 2016

Andrzej Czarnecki  University of Alberta
with R. Szafron, M. Dowling, J. Piclum

Outline

Muon decay in orbit

- two energy regions
- different approaches to radiative corrections



g -factor of a bound electron

- towards $(Z\alpha)^5$ effects

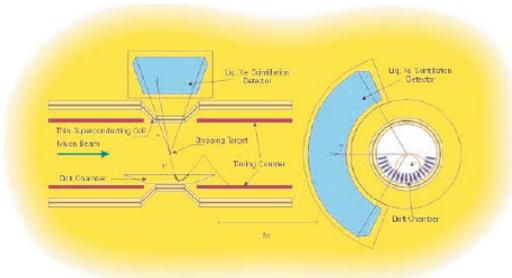
Lepton flavor violation: $\mu \rightarrow e\gamma$

Previous experiment (MEGA @ Los Alamos): $\text{BR}(\mu \rightarrow e\gamma) < 10^{-11}$

New bound (MEG @ Paul Scherrer Institute)

$< 4.2 \cdot 10^{-13}$

arXiv:1605.05081



This corresponds to the transition dipole moment

$$d_{\mu \rightarrow e} \lesssim 3.5 \cdot 10^{-27} e \cdot \text{cm}$$

Sensitive to
the heaviest
"new physics"

For comparison: electron EDM $d_e < 0.87 \cdot 10^{-28} e \cdot \text{cm}$

10.1126/science.1248213

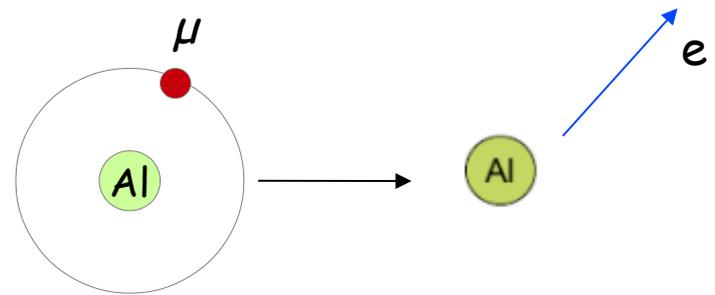
muon $g-2$

$$d_\mu < 3 \cdot 10^{-22} e \cdot \text{cm}$$

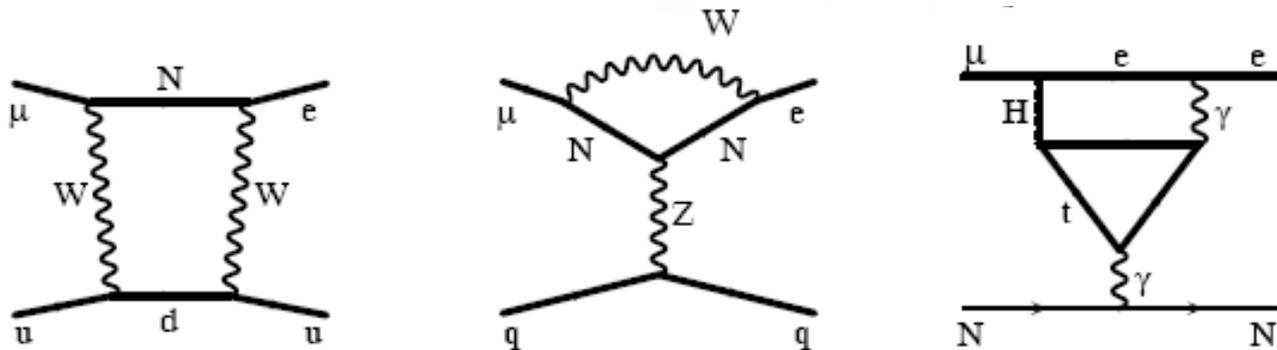
Muon-electron conversion: probes various types of interactions

Non-dipole interactions are not (directly) probed by processes with external photons, by gauge invariance requirements.

New process: muon-electron conversion
(as well as $\mu \rightarrow eee$)

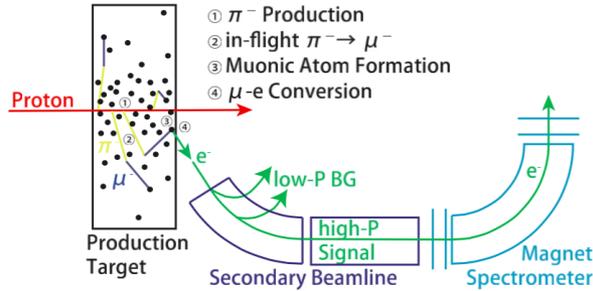


Variety of mechanisms:



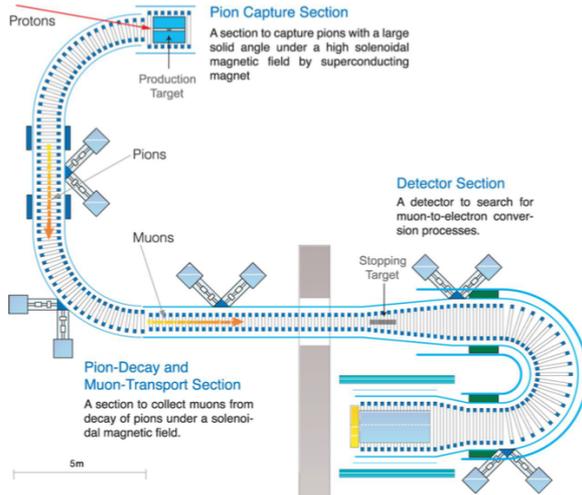
Muon-electron conversion plans (The Next Big Thing in muon physics)

DeeMe
J-PARC



starts 2016;
aims for $1e-13$ (graphite target),
followed by $1e-14$ (SiC target)

COMET
Phase 1
J-PARC

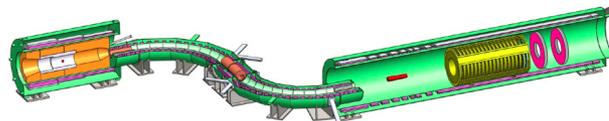


$7e-15$

COMET
Phase 2
J-PARC

$2.6e-17$

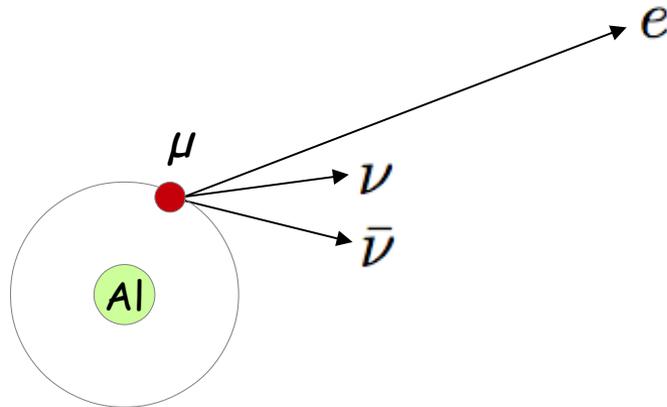
Mu2e
Fermilab



$2e-17$

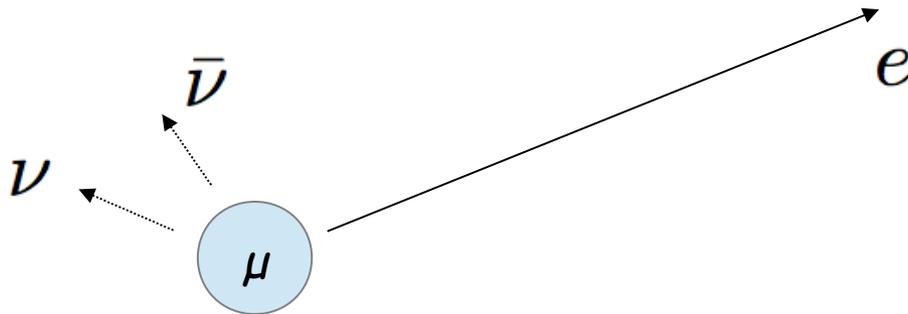
Background for the conversion search

Normal decay of the muon bound in the atom can produce high-energy electron,

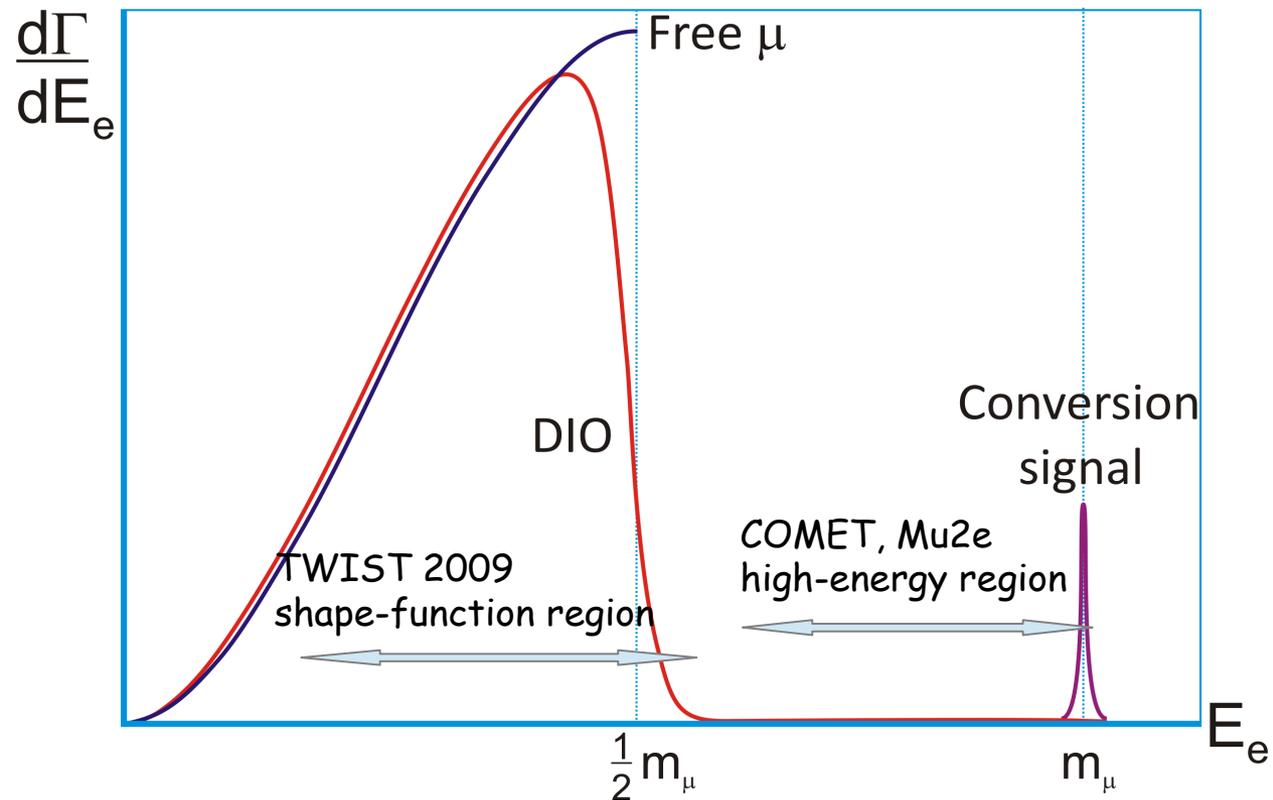


Spectrum has to be well understood.

Electron spectrum in a bound muon decay



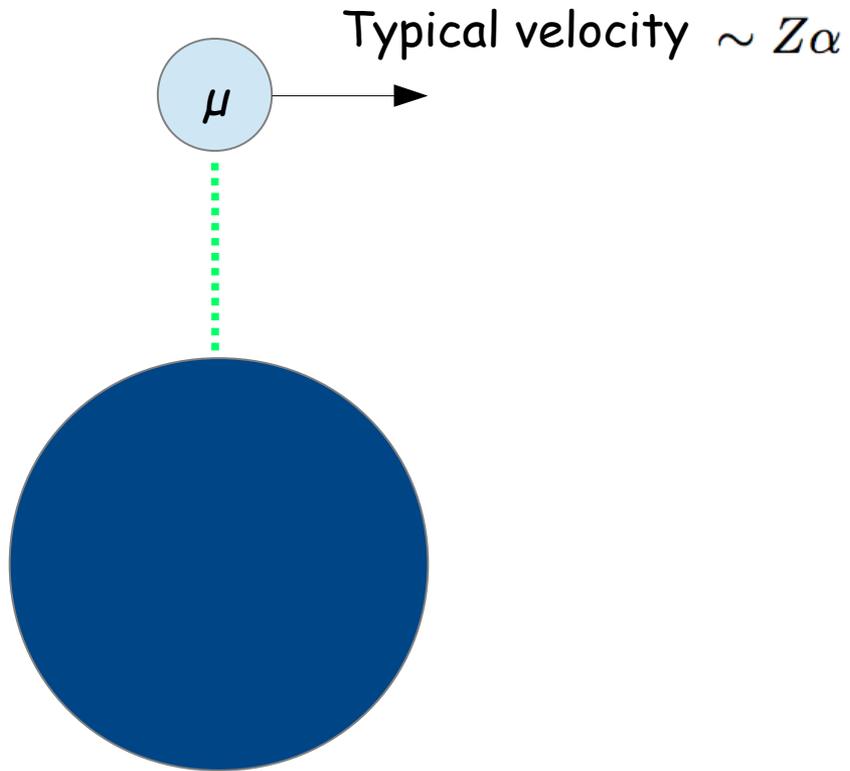
Electron energy can be as large as the whole muon mass



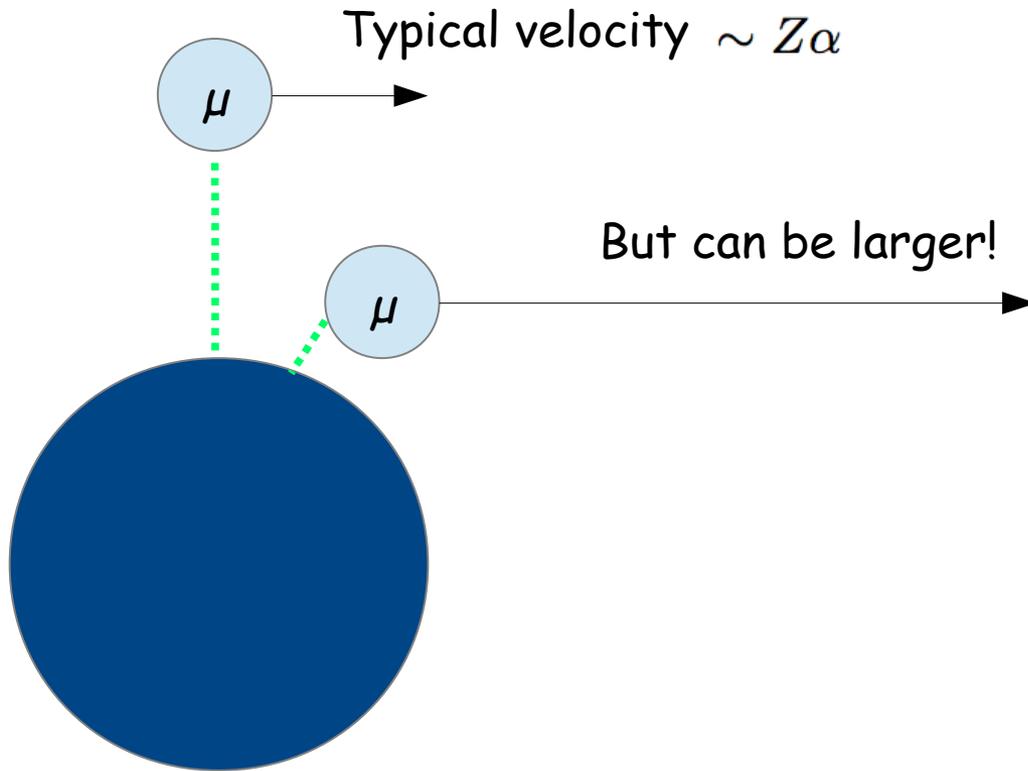
**Muon decay-in-orbit spectrum:
the shape-function region**

Experiment: TWIST

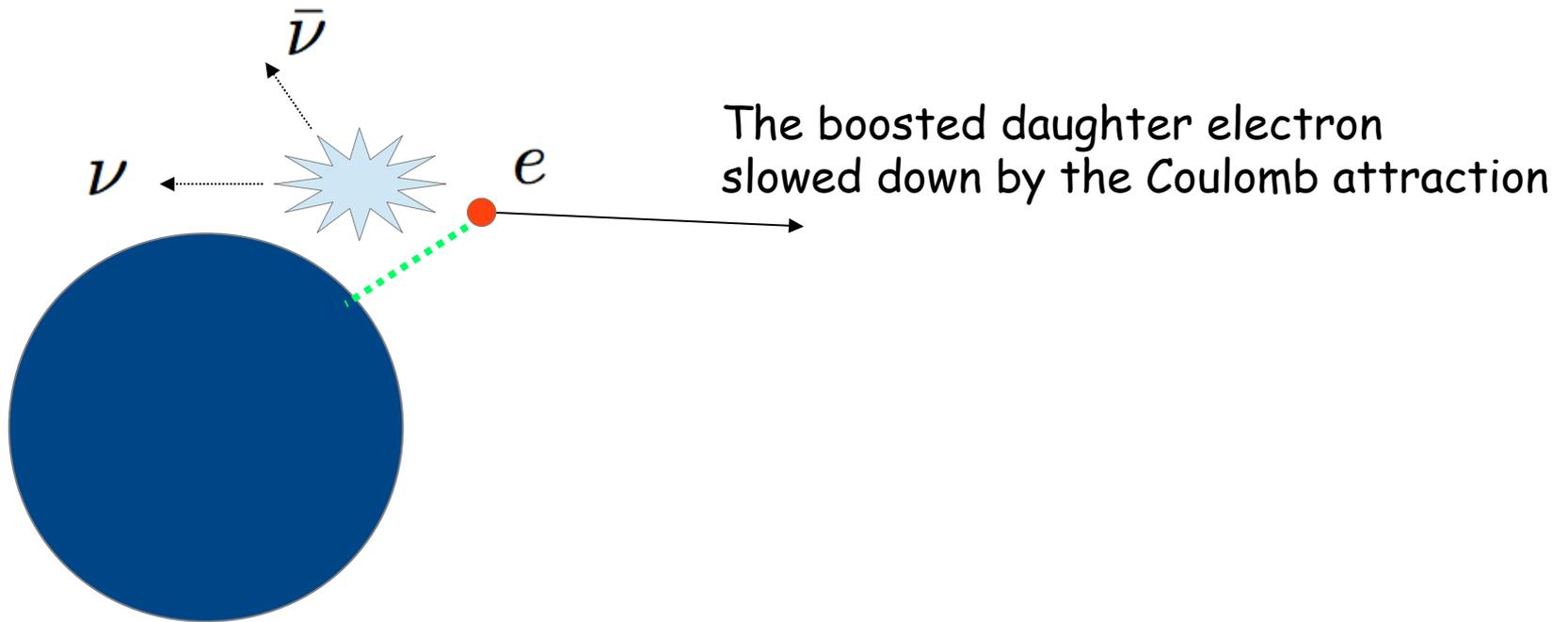
Two effects: muon motion & Coulomb attraction



Two effects: muon motion & Coulomb attraction



Two effects: muon motion & Coulomb attraction



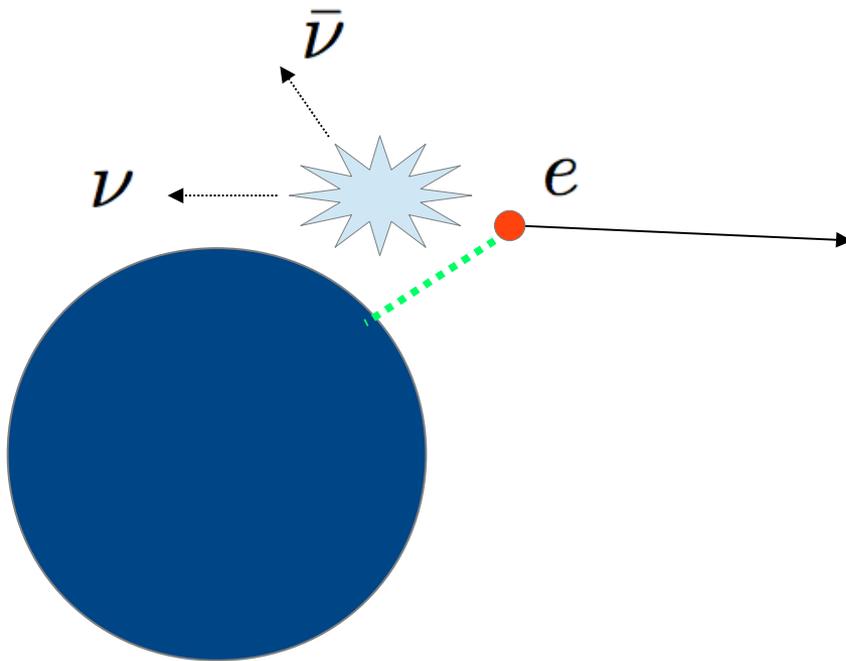
Two effects: muon motion & Coulomb attraction

Net effect:

- In the decay rate: almost none;
only time dilation

$$\Gamma \rightarrow \left(1 - \frac{(Z\alpha)^2}{2}\right) \Gamma$$

Überall, Phys. Rev. 119, 365 (1960)



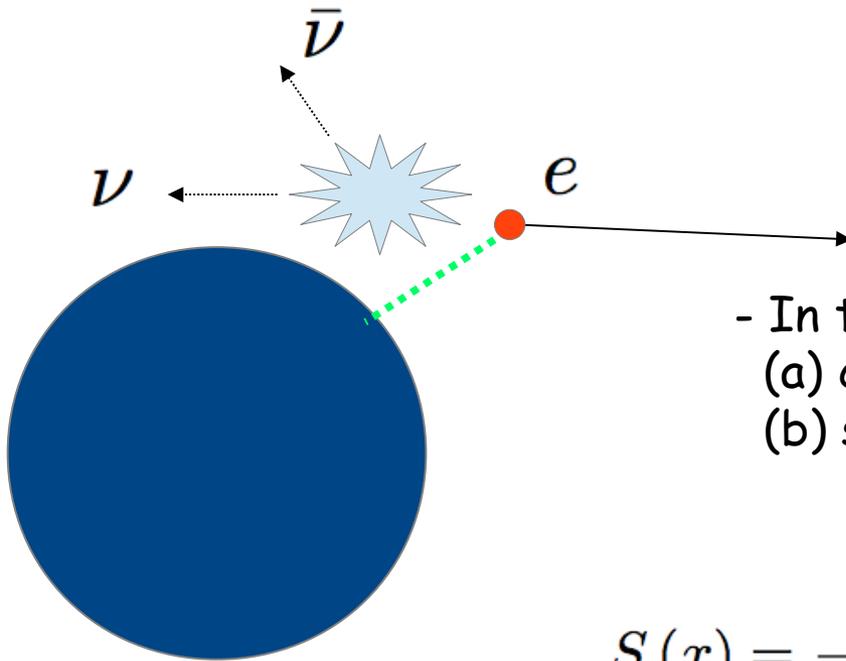
Two effects: muon motion & Coulomb attraction

Szafron, AC, PRD 92 (2015) 053004

Net effect:

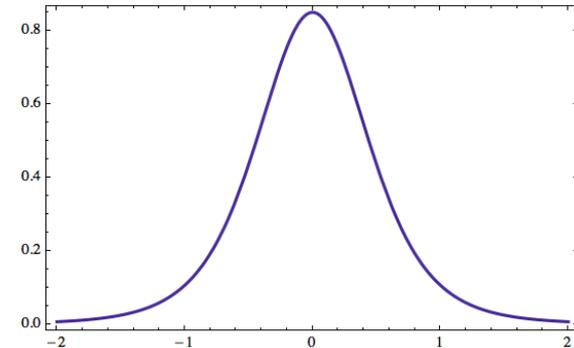
- In the decay rate: almost none; only time dilation

$$\Gamma \rightarrow \left(1 - \frac{(Z\alpha)^2}{2}\right) \Gamma$$



- In the electron energy spectrum:
 - (a) computable shift
 - (b) smearing → "shape function"

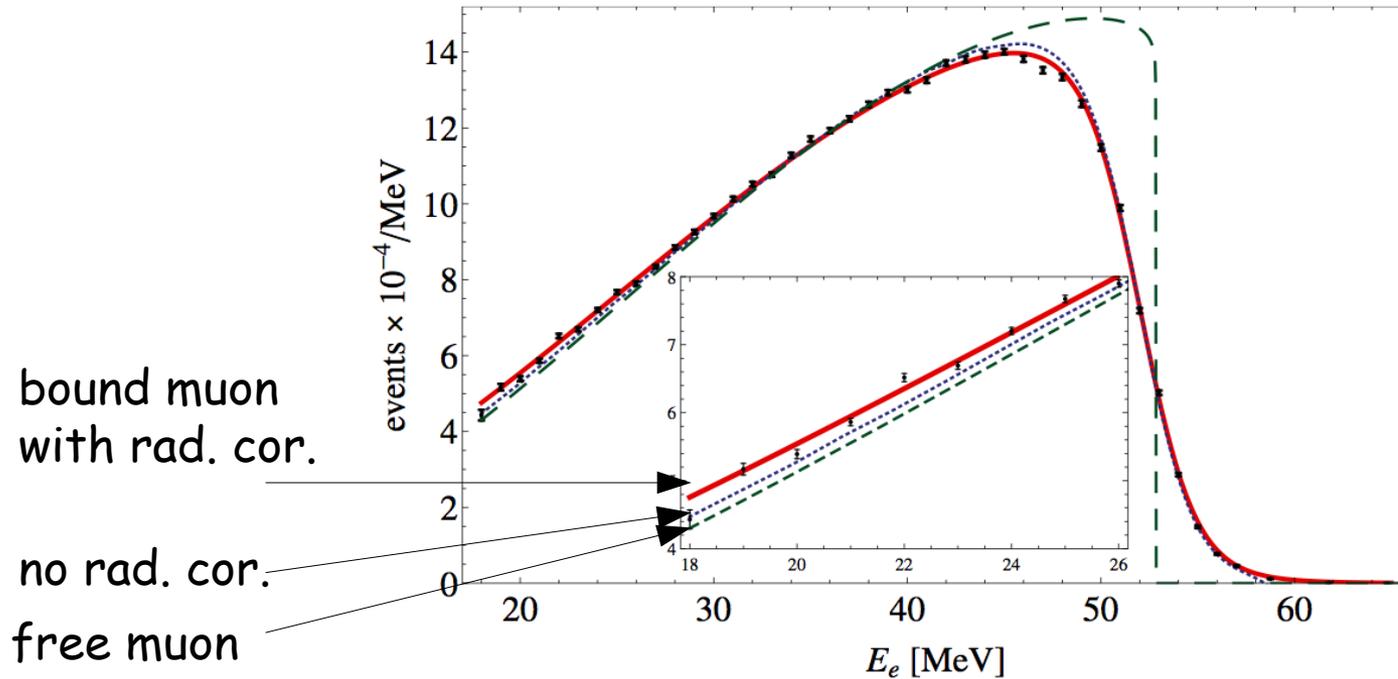
$$S(x) = \frac{8}{3\pi [1 + x^2]^3}$$



Previously used in heavy mesons, where it cannot be computed from first principles, but can be experimentally accessed.

Mannel, Neubert,
Bigi, Shifman, Uraltsev, Vainshtein

Comparison with measurement: TWIST

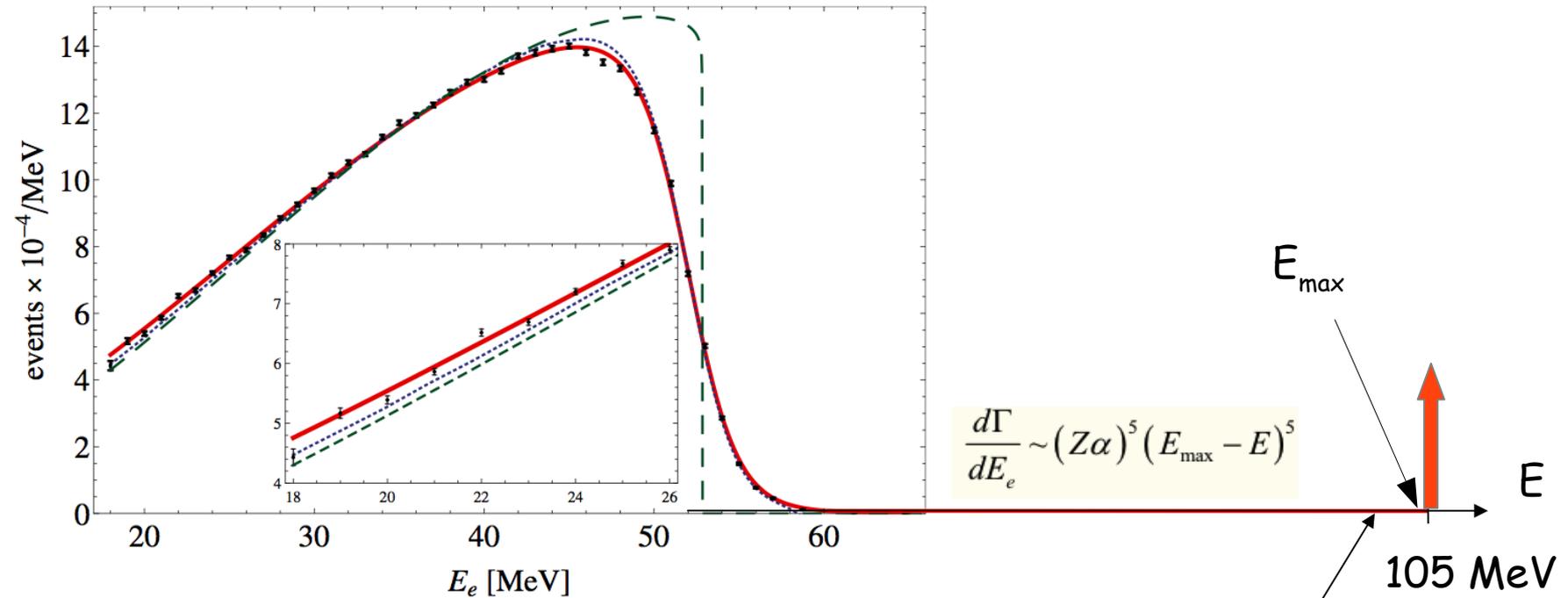


The spectrum is modified very significantly: effects $\sim 1/Z\alpha$

**Muon decay-in-orbit spectrum:
the high-energy region**

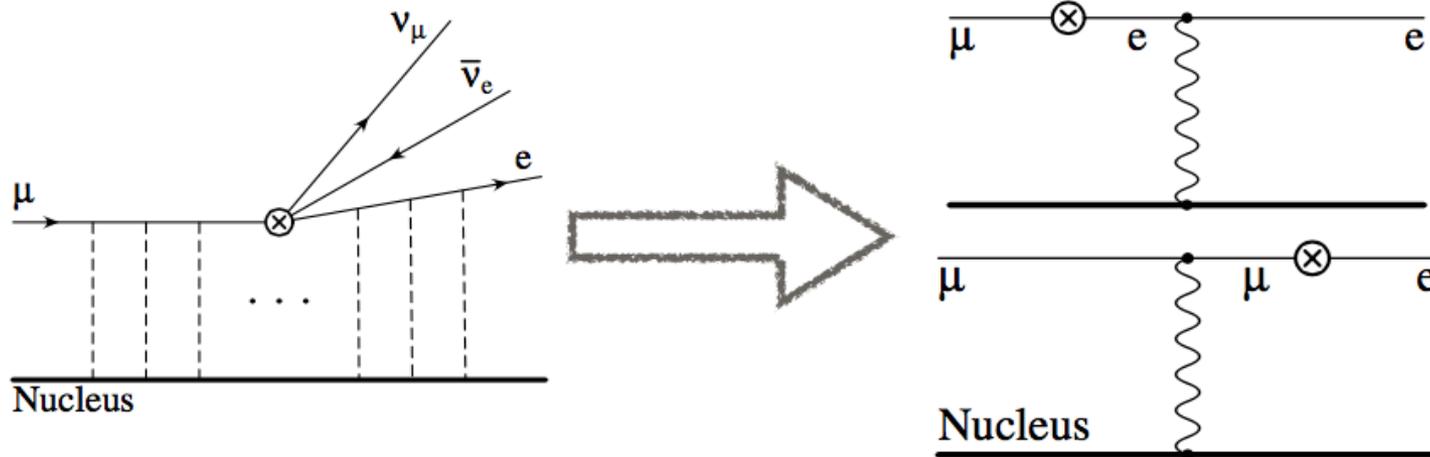
Experiments: Mu2e and COMET

Spectrum of the bound muon decay



Main background for the conversion signal

Origin of the $(E_{\max} - E)^5$ suppression



Neutrinos get no energy;

The nucleus balances electron's momentum, takes no energy.

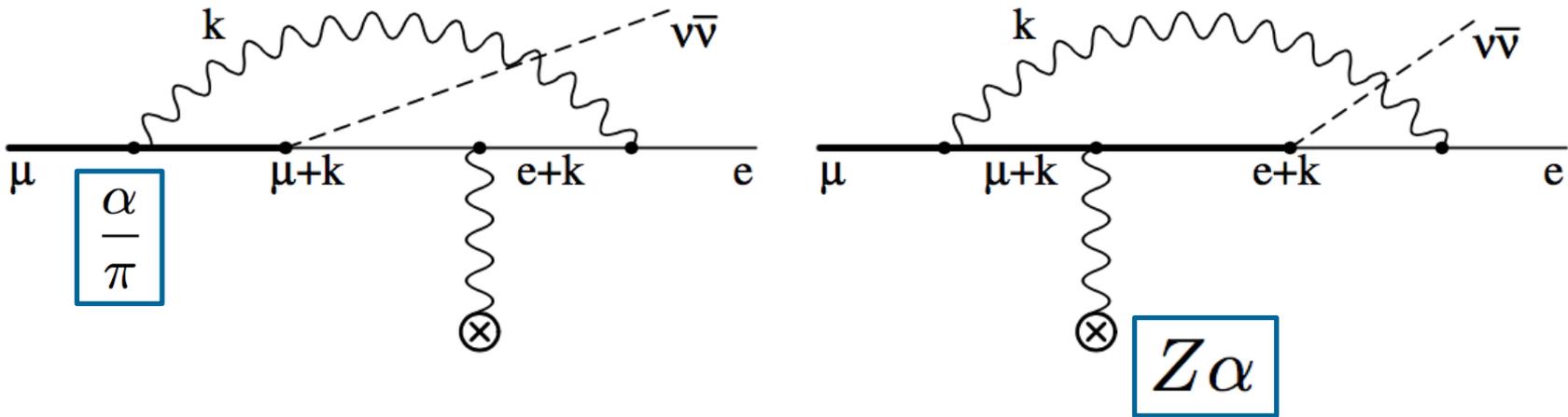
Near the end point:

Radiative corrections to the electron spectrum

Expansion near the end-point $\frac{m_\mu}{\Gamma_0} \frac{d\Gamma}{dE} = \sum_{ijk} B_{ijk} \Delta^i (\pi Z\alpha)^j \left(\frac{\alpha}{\pi}\right)^k$

Three "small" parameters:

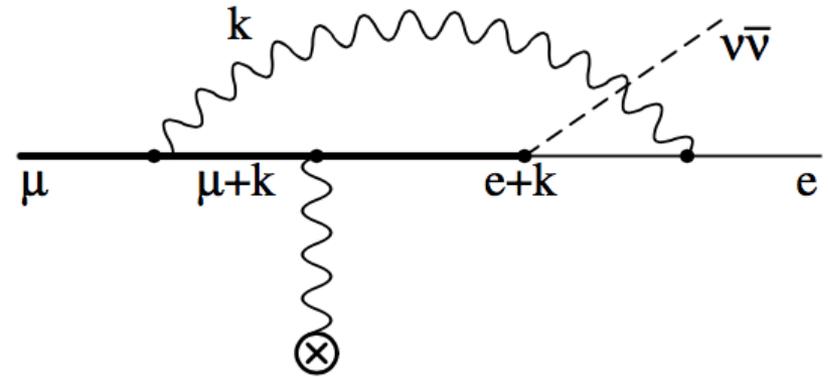
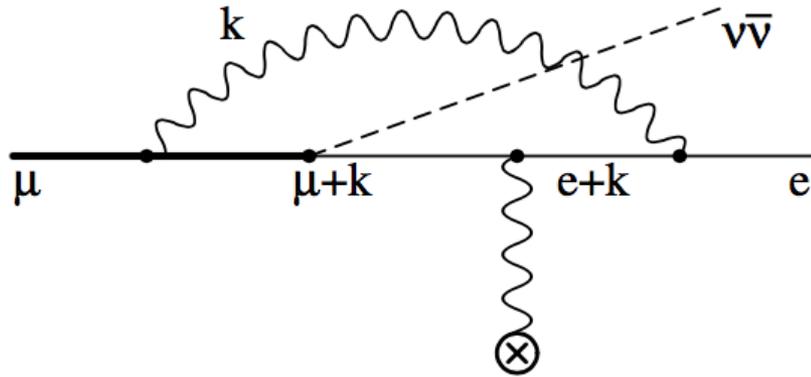
$$\Delta = \frac{E_{\max} - E}{m_\mu}$$



The expansion starts with B_{550}

The first radiative correction is B_{551}

Radiative corrections to the electron spectrum



Expansion near the end-point $\frac{m_\mu}{\Gamma_0} \frac{d\Gamma}{dE} = \sum_{ijk} B_{ijk} \Delta^i (\pi Z\alpha)^j \left(\frac{\alpha}{\pi}\right)^k$

Competing effects:

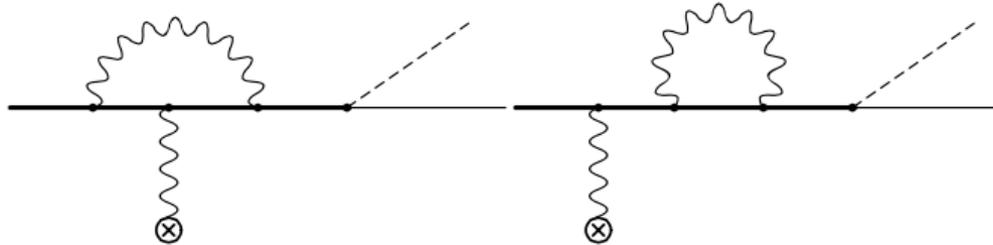
- vacuum polarization in the hard photon; and
- self-energy and real radiation

$$\frac{B_{551}}{B_{550}} = \delta_H + \delta_S \ln \Delta$$

$$\delta_H = 6.31 - \frac{26}{15} \ln \frac{m_\mu}{m_e}$$

$$\delta_S = 2 \ln 2 - 2 + 2 \ln \frac{m_\mu}{m_e}$$

Results from the first radiative corrections



The most important correction is due to soft photons.
Can be summed up to all orders (exponentiated).

$$B_{550} + \frac{\alpha}{\pi} B_{551} \rightarrow B_{550} \left[\Delta \frac{\alpha}{\pi} \delta_S + \frac{\alpha}{\pi} \delta_H \right]$$

$$\delta_S = 10.1$$

number of electrons in the
end-point bin of 1 (0.1) MeV is reduced by 11% (16%)

Magnetic moment of a bound electron

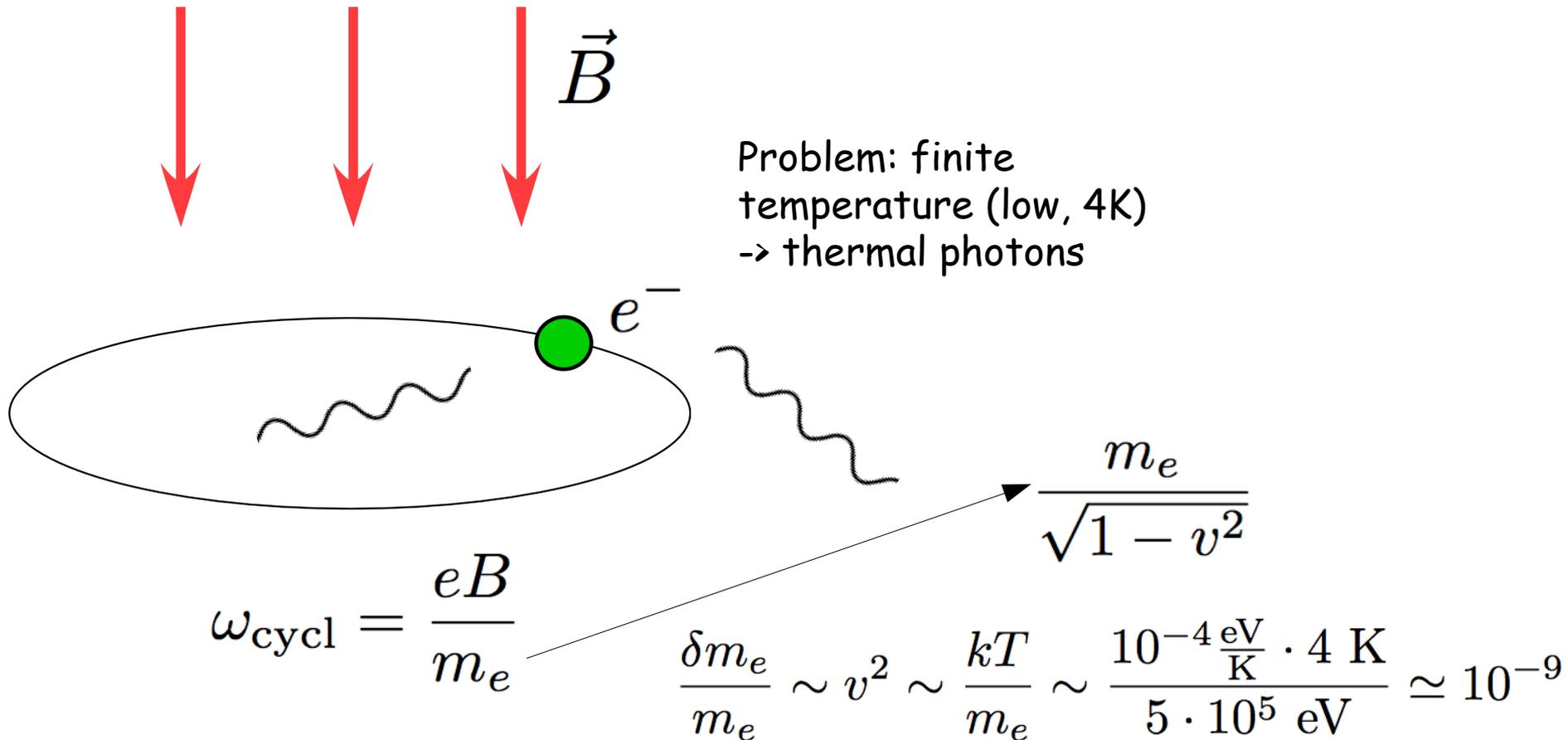
Why useful?

- determination of the electron mass
- future determination of α
- indirectly related to muon $g-2$ (muonium)

Why interesting?

- quantum effects in external field
- simple system, model for more complex ones
- numerical estimates exist for large Z
- should be analytically feasible for small Z

Determination of the electron mass (20th century)

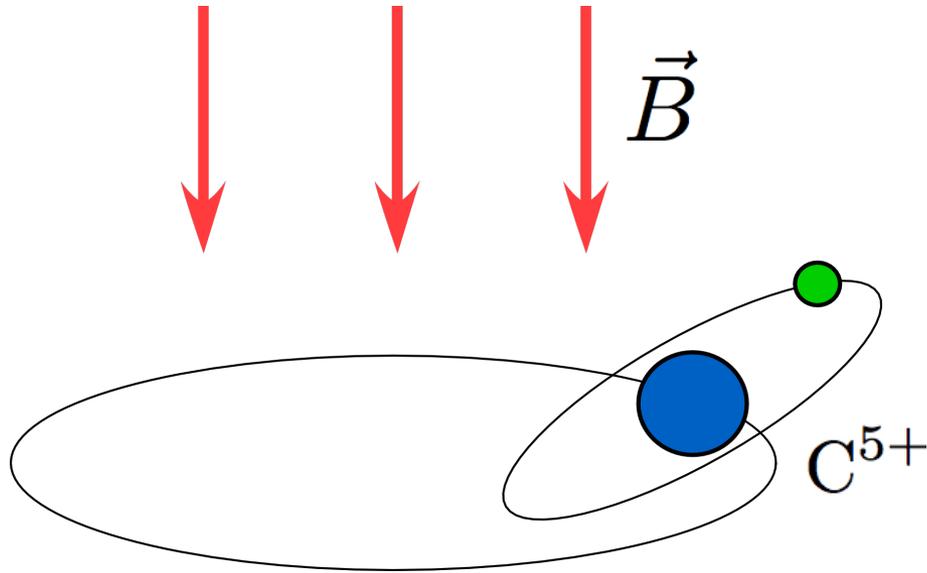


Determination of the Electron's Atomic Mass and the Proton/Electron Mass Ratio via Penning Trap Mass Spectroscopy

$$\frac{m_e}{u} = 0.000\,548\,579\,911(1)$$

2ppb relative error

21st century: anchor the electron in an ion

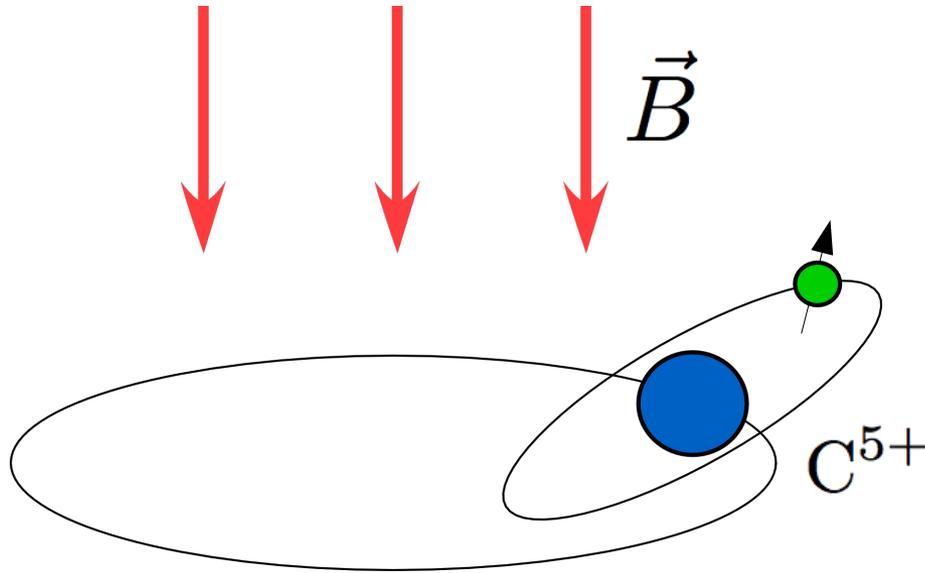


The speed of the electron becomes even higher,

$$v_e^2 \sim (Z\alpha)^2 \sim 10^{-3}$$

but is calculable
in bound-state QED.

21st century: anchor the electron in an ion



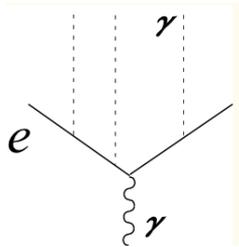
Larmor frequency $\omega_L = \frac{geB}{2m_e}$

Cyclotron frequency $\omega_{cycl} = \frac{ZeB}{M}$

$$m_e = \frac{g}{2Z} \frac{\omega_{cycl}}{\omega_L} M$$

Bound-electron $g-2$: the leading effect

Breit 1928: energy correction due to magnetic field in the hydrogen ground state.

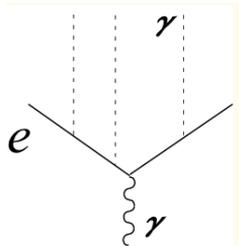


$$\delta E = e \int d^3x f^2 v^* [1 - i\gamma \boldsymbol{\Sigma} \cdot \hat{\mathbf{r}} \gamma^5] \gamma^5 \mathbf{A} \cdot \boldsymbol{\Sigma} [1 + i\gamma \boldsymbol{\Sigma} \cdot \hat{\mathbf{r}} \gamma^5] v$$

$$g = 2 \cdot \frac{1}{3} \left(1 + 2\sqrt{1 - (Z\alpha)^2} \right) \simeq 2 \left(1 - \frac{(Z\alpha)^2}{3} \right)$$

Bound-electron g-2: the leading effect

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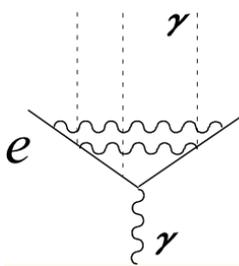
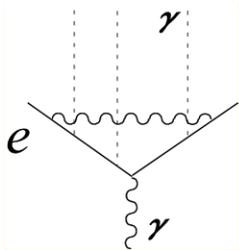
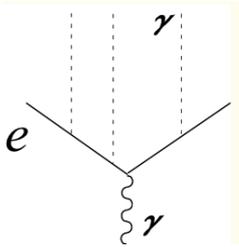
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$$g = 2 \cdot \frac{1}{3} \left(1 + 2\sqrt{1 - (Z\alpha)^2} \right) \simeq 2 \left(1 - \frac{(Z\alpha)^2}{3} \right)$$

Important: dependence on alpha; may be exploited to determine its value.
(Use ions with various Z)

Shabaev, Glazov, Oreshkina, Volotka, Plunien, Kluge, Quint Yerokhin, Berseneva, Harman, Tupitsyn, Keitel: PRL (2016)

Bound-electron $g-2$: binding and loops



$$g = 2 - \frac{2(Z\alpha)^2}{3} - \frac{(Z\alpha)^4}{6} + \dots$$

$$+ \frac{\alpha}{\pi} \left[1 + \frac{(Z\alpha)^2}{6} + (Z\alpha)^4 (a_{41} \ln Z\alpha + a_{40}) + \dots \right]$$

$$+ \underbrace{\left(\frac{\alpha}{\pi} \right)^2 \left[-0.65.. \left(1 + \frac{(Z\alpha)^2}{6} \right) + (Z\alpha)^4 (b_{41} \ln Z\alpha + b_{40}) + \dots \right]}_{\text{two-loop corrections}}$$

$$b_{41} = \frac{28}{9}$$

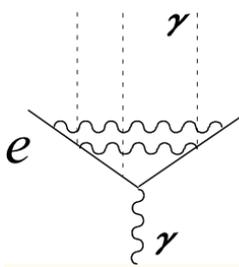
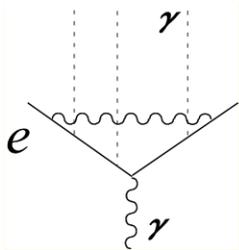
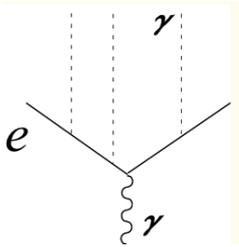
$$b_{40} = -16.4$$

Pachucki,
AC
Jentschura,
Yerokhin
2005

Together with experiments (Wolfgang Quint et al),
this improved the accuracy of m_e by about a factor 3,

$$\frac{m_e}{u} = 0.000\,548\,579\,909\,32(29)(1)$$

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$$\frac{m_e}{u} = 0.000\ 548\ 579\ 909\ 067\ (17)$$

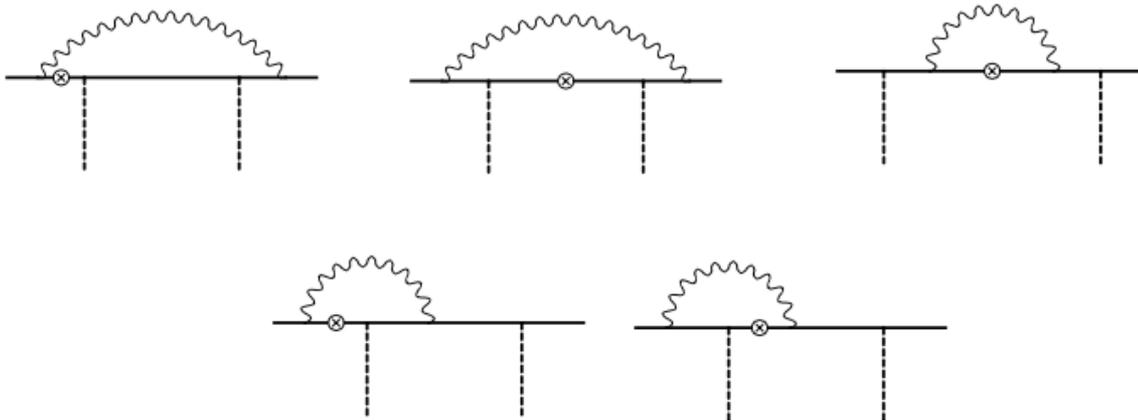
Nature 2014
Sturm et al

Next theory challenge:
 $(Z\alpha)^5$ effects.

Numerical result for
 $\alpha(Z\alpha)^5$ but not $\alpha^2(Z\alpha)^5$

Yerokhin, Indelicato, Shabaev

Towards an analytic result for $a(Z\alpha)^5$



The result is gauge-invariant; but not yet complete.

What if the magnetic field couples to an external line?

Magnetic correction to the wave function

$$\Delta g = \frac{2|e|\hbar}{\mu_0 m B} \sum_n^{n \neq a} \frac{\langle a | \delta U | n \rangle \langle n | \vec{\alpha} \cdot \vec{A} | a \rangle}{E_a - E_n}$$

This sum can be done exactly with help of virial identities

S. Karshenboim, V. Ivanov, and V. Shabaev

$$\Delta g = \frac{2\kappa m}{j(j+1)} \langle n\kappa | \delta U \frac{\kappa}{m^2} (I - |n\kappa\rangle \langle n\kappa|) \left[\left(E_{n\kappa} - \frac{m}{2\kappa} \right) r i \sigma_y + m r \sigma_x + \alpha Z i \sigma_y - \kappa \sigma_z \right] |n\kappa\rangle$$

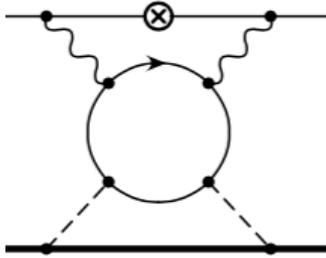
The result is especially simple if the short-distance perturbation is energy-independent (like the Uehling potential)

$$\Delta g = \frac{16}{3} \left(1 - \frac{1 - \sqrt{1 - (Z\alpha)^2}}{2(Z\alpha)^2} \right) \langle \delta^3(r) \rangle_\psi \simeq 4 \langle \delta^3(r) \rangle_\psi$$

This is just Lamb

Next goal: $\alpha^2(Z\alpha)^5$ corrections to g

Example:

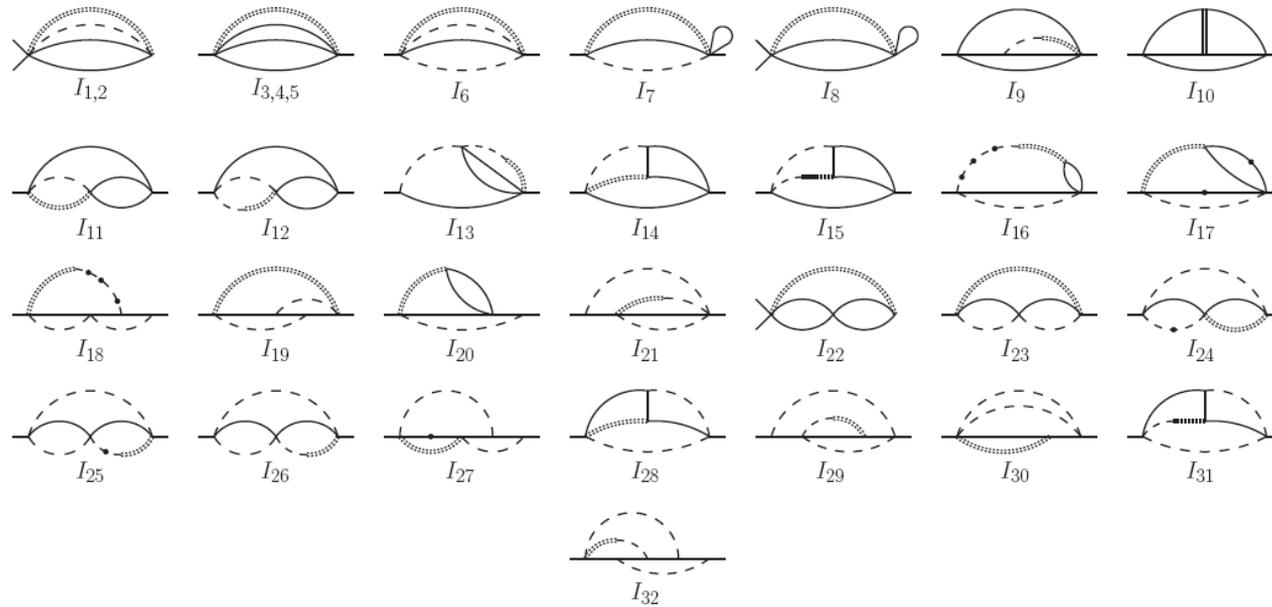


Related diagrams:

Karshenboim and Milstein, 2000-2002

More than 300 contributions.

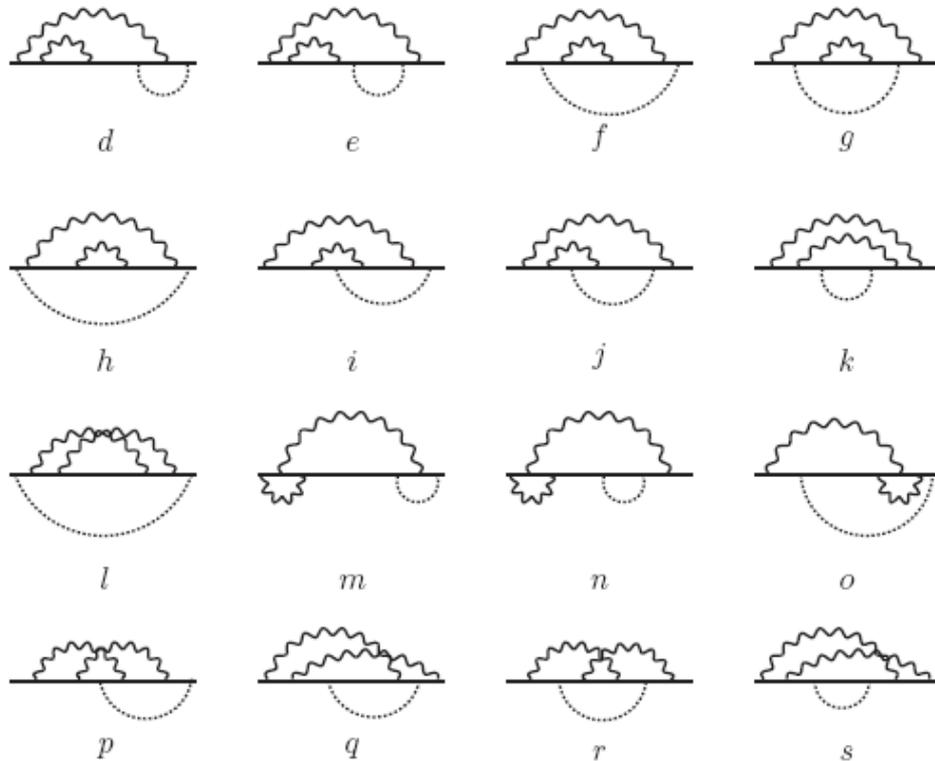
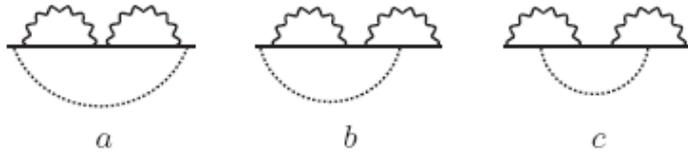
A set of 32 master integrals



Typical expression

$$I_{24} = G(0, 1, 2, 1, 0, 1, 0) = \frac{2\pi^2}{\epsilon} - 162.745878930257(1) + 640.681562239(2)\epsilon - 9490.745115169417(3)\epsilon^2 + \mathcal{O}(\epsilon^3),$$

Reevaluation of the $\alpha^2(Z\alpha)^5$ Lamb shift



$$\delta E_{a-s} = \frac{\alpha^2(Z\alpha)^5}{\pi n^3} \left(\frac{\mu}{m}\right)^3 m [-7.72381(4)]$$

Dowling, Mondejar, Piclum, AC, PRA 81, 022509

Previous results

-7.61(16) Pachucki 1994

-7.724(1) Eides and Shelyuto, 1995

Summary

- * Binding modifies the muon decay and the electron g -factor
- * Theory of both effects is more fun than for free particles
- * Synergy with beautiful experiments: lepton-flavor violation, mass of the electron and, in future, the fine structure constant.
- * For g : $\alpha(Z\alpha)^5$ effects almost finished; $\alpha^2(Z\alpha)^5$ hopefully soon.
- * Opportunities for more theoretical improvement...