

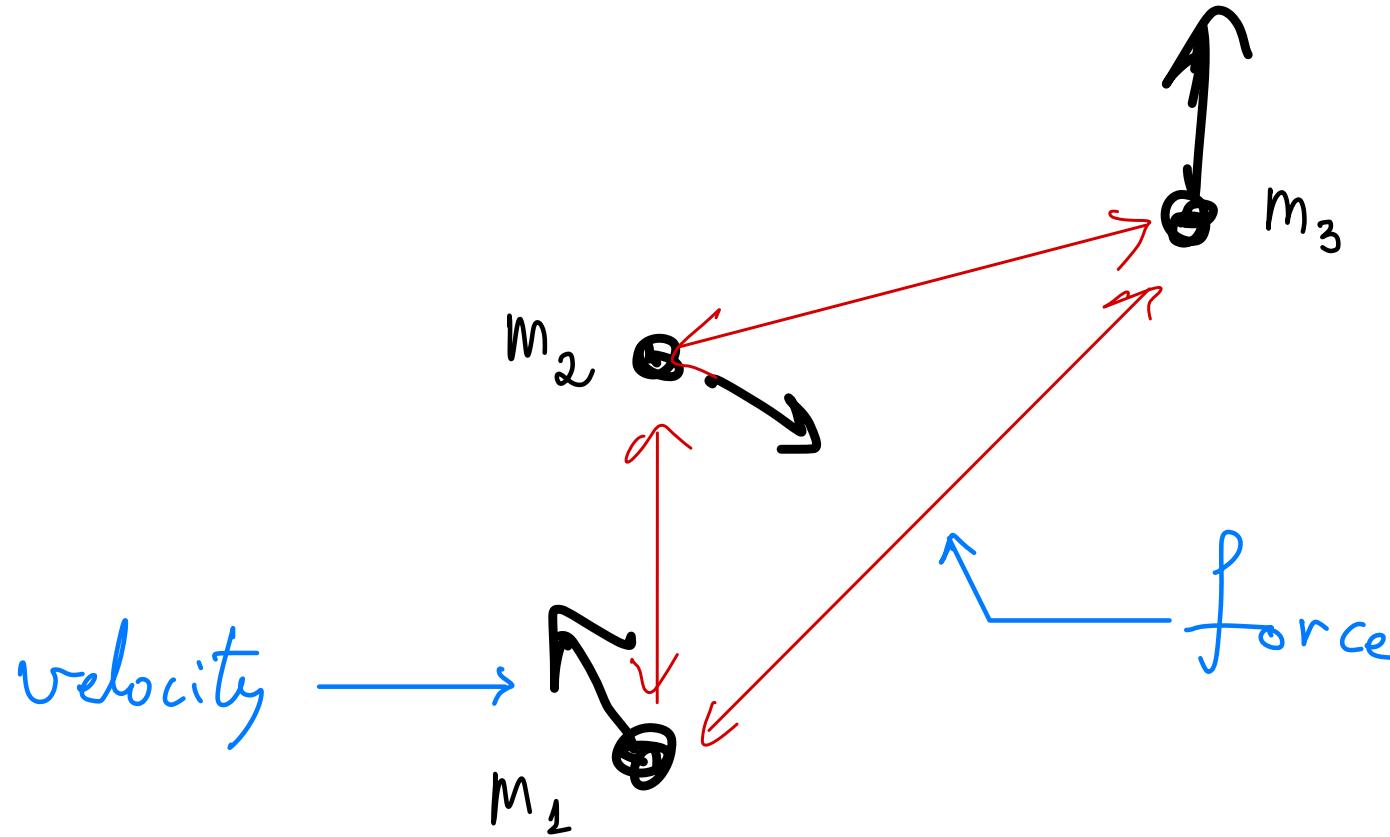
Natural dynamical reduction of the three-body problem

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based on arXiv:2107.12372

- * The problem
- * Coord. for center of mass frame - $\vec{w} \in \mathbb{F}^3$
- * Orientation - geometry decomposition -
formulation w. $\vec{\mathcal{T}}$
- * Applications
- * Perspective

The three-body problem



e.g. $F = \frac{G m_1 m_2}{r^2}$

$$\mathcal{L} = \frac{1}{2} \sum m_a \dot{r}_a^2 - \left(\frac{G m_1 m_2}{r_{12}} + \text{c.c.} \right)$$

General background

Newton explained planetary motion.

Universal law of gravitation

2-body problem.

What about 3-body ?

General background

The moon, perturbation theory, central configurations,
Poisson bracket, secular approximation,
chaos, modern topology,
computer simulations, statistical theory

Newton, d'Alembert, Clairaut, Euler, Lagrange, Laplace,
Gauss, Jacobi, Hill, Poincaré,
Levi-Civita, Lemaître; Hylleraas, Fock,
Agekian & Anosova, Valtonen, Monaghan, Heggie, Hut, McMillan,
Smale, Saari, C. Moore, Chenciner, Montgomery & Moeckel
Naoz, Peretz, Stone.

Dynamical reduction

- * Simplify the dynamics through choice of variables
- * Conserved quantities enable it
- * Examples
 - Small oscillations w. $n > 1$ d.o.f.
 - Central force & effective radial potential

My road

- * Stone & Leigh (2019) - Approx. Stat. soln in closed form
- * BK (2020) - Flux-based theory to appear w. collab.*
Regularized phase space volume

$$\mathcal{T}(E_0, \vec{L}_0) = \int \prod_{Q=1}^3 \left(d^3 r_Q d^3 p_Q \right) S(E_0, \vec{L}_0, \vec{P}, \vec{R}_{cm})$$

see also Ginat & Perets 2020

* Dandekar
Mazumdar
Lederer

Coordinates for center of mass frame

$$\vec{r}_1, \vec{r}_2, \vec{r}_3$$

3 vector d.o.f.

$$\vec{0} = m_1 \vec{r}_1 + m_2 \vec{r}_2 + m_3 \vec{r}_3$$

1 vector constraint

In 2-body

$$\vec{r} = \vec{r}_1 - \vec{r}_2$$

In 3-body relative positions (transl. invar.)

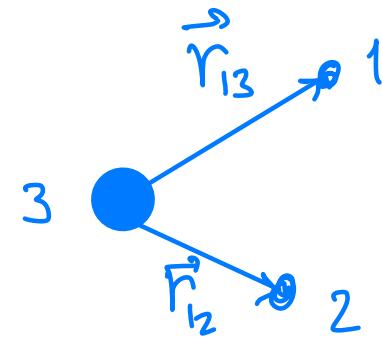
$$\vec{r}_{ab} = \vec{r}_a - \vec{r}_b : \vec{r}_{12}, \vec{r}_{23}, \vec{r}_{31}$$

$$\text{not independent} : \vec{r}_{12} + \vec{r}_{23} + \vec{r}_{31} = 0$$

Popular choices of coord.

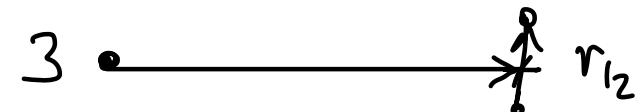
* Planetary $M_3 \gg m_1, m_2$

$$\vec{r}_{13}, \vec{r}_{23}$$



* Lunar / hierarchical

$$\vec{r}_{12}, \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2} - \vec{r}_3$$



Inspired by Lagranges solution to the cubic (1771)

$$Qx^3 + bx^2 + cx + d = 0$$

Define

$$\vec{w} = \vec{r}_1 + \eta \vec{r}_2 + \bar{\eta} \vec{r}_3$$

$$\eta = \exp(2\pi j/3)$$

Translation invariant.

Transforms nicely under perm.: $1 \rightarrow 2 \rightarrow 3 \rightarrow 1$
 $2 \leftrightarrow 3$

$$w \rightarrow \bar{\eta} w$$
$$w \rightarrow \frac{1}{w}$$

Toward formulation

- (a) translations ✓ \vec{w}
- (b) rotations $\vec{w} \rightarrow \vec{\omega}, w$
- (c) rotations within plane, $\mathbb{C}^2/\text{U}(1)$, $w \rightarrow \vec{g}$
- (d) Transform $\vec{\omega} \rightarrow \vec{\tau}$

Appendices

- * Non coord. velocities & rigid body
Euler eq's from L, H
- * Bi-complex numbers & isotropic (2d) oscillator

Orientation — Geometry decomposition

$$\vec{J}^2 = \text{const}$$

chaotic core

Geometry space — \vec{g}
 $\leftrightarrow (r, \theta, \varphi)$
pipe joint

Formulation

$$L_J(\vec{J}, \vec{g}) := L - \vec{\omega} \cdot \vec{J} = L_0 + L_1 + L_2$$

↑
 geometry space motion potential

$$L_0 = \frac{1}{8I} \dot{I}^2 + \frac{3M_3}{8M} \frac{r^4}{I} (\dot{\theta}^2 + \sin^2\theta \dot{\phi}^2) - \left(\frac{GM_2M_3}{r_{23}} + c.c. \right)$$

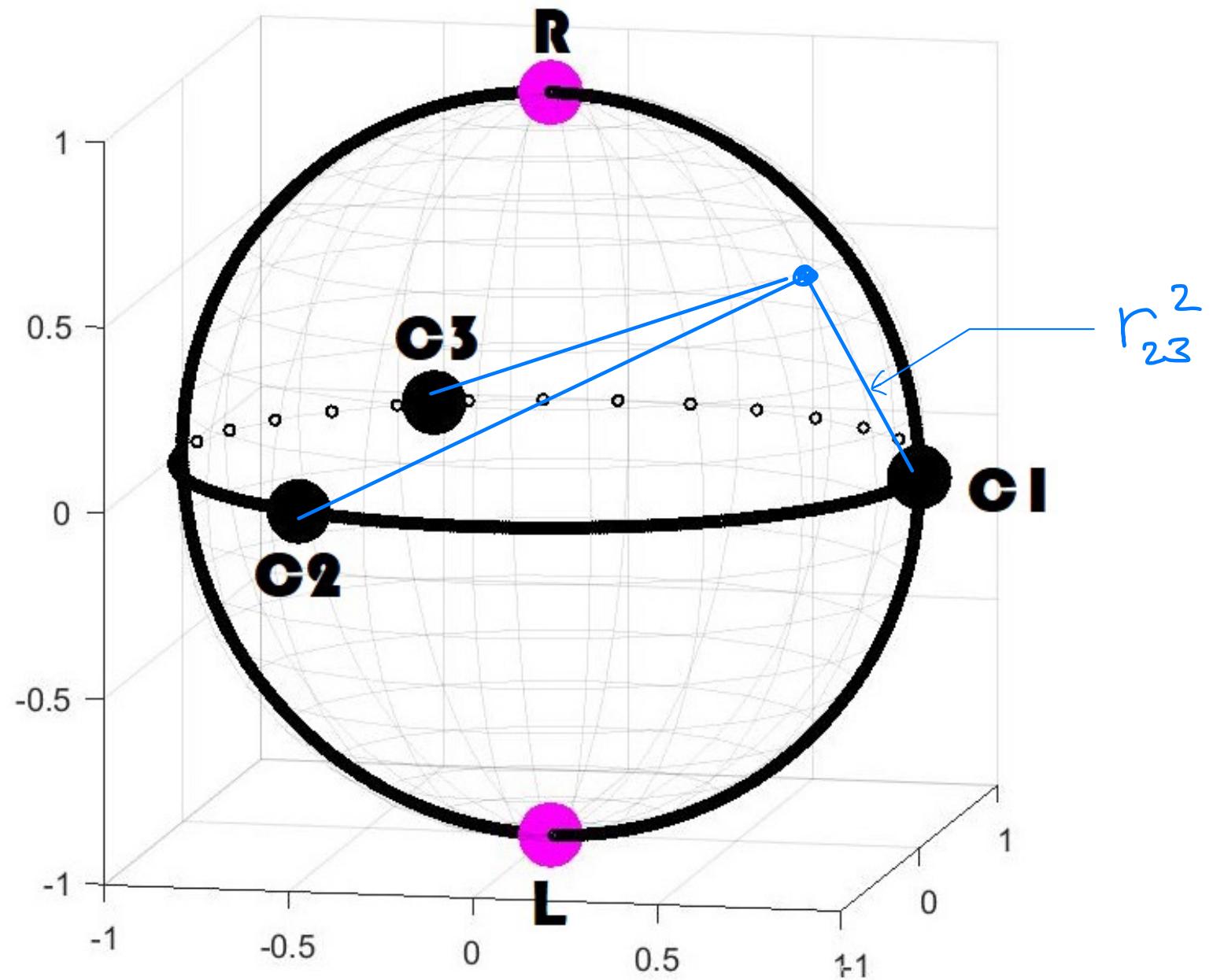
$$I = \sum m_a r_a^2 \uparrow \frac{1}{M} (m_2 m_3 r_{23}^2 + \text{cyclic}) \quad M = M_1 + M_2 + M_3$$

$$r^2 = r^2(1 \sin \theta \cos \psi)$$

$$M_3 = m_1 m_2 m_3$$

$$r_{31}^2 = " (" - " \cos(\varphi - \frac{2\pi}{3}))$$

Shape sphere metric conformal to round S^2 .



$$L_1 = J_3 \frac{L_G}{I} ; \quad L_G = -\frac{r^2}{2M} \left[m_2 m_3 ((1 - \cos \theta - \sin \theta \cos \varphi) \dot{\phi} - \sin \varphi \dot{\theta}) + \text{cyc.} \right]$$

↑
in "r₊-gauge"

Magnetic-like 2-form

$$B = \frac{3}{2} M m_3 J_3 \frac{\sin \theta d\varphi d\theta}{[m_2 m_3 (1 - \sin \theta \cos \varphi) + \text{cyc.}]^2}$$

monopole charge at $\vec{q} = 0$.

$$-L_2 = \frac{1}{2I} J_3^2 + \frac{M}{2m_3} \frac{1}{\Delta^2} \bar{I}^{\alpha\beta} J_\alpha J_\beta \quad \alpha, \beta = 1, 2$$

$$\Delta^2 = \frac{3}{4} r^4 \cos^2 \theta \quad \Delta \text{ is twice triangle area}$$

$$\bar{I}^{\alpha\beta} = \frac{1}{M} [m_2 m_3 \rho_{23}^\alpha \rho_{23}^\beta + \text{cyc.}] \quad \sim I^{-1}$$

$$\rho_{23} = ir \left[\cos \frac{\theta}{2} - e^{-i\phi} \sin \frac{\theta}{2} \right]$$

in τ_+ gauge

$$\{J_i, J_j\} = -\epsilon_{ijk} J_k$$

centrifugal potential

Equations of motion

$$\frac{d}{dt} \tau_i = - \{ \tau_i, L_J \}$$

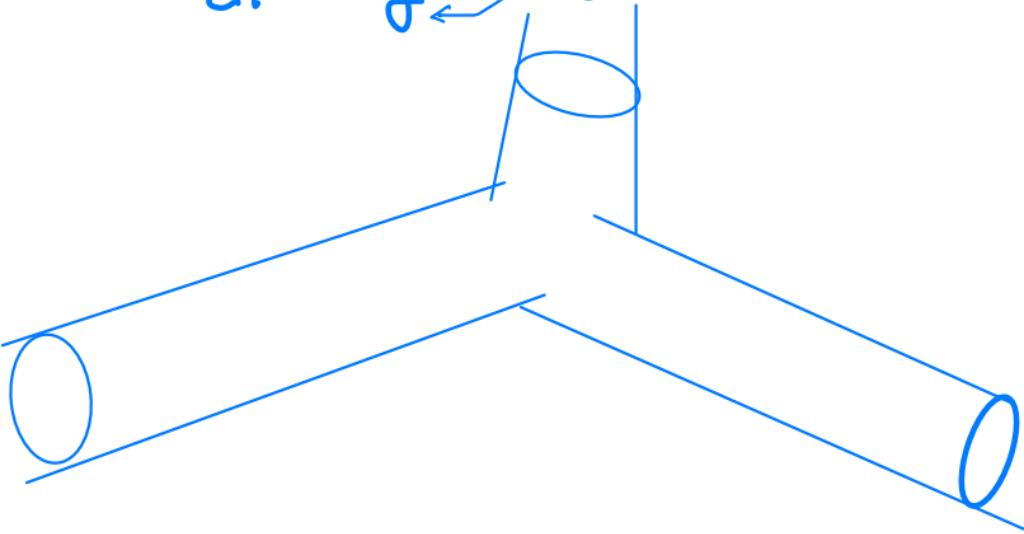
$$\frac{d}{dt} \left(\frac{\partial L_J}{\partial \dot{q}^s} \right) = \frac{\partial L_J}{\partial q^s} \quad q^s = (r, \theta, \varphi)$$

Applications

- * to statistical soln
- * to exact solutions

Regularized phase space volume

$$d\Gamma \sim \frac{dF}{\Omega} \quad \text{reg. phase vol.}$$

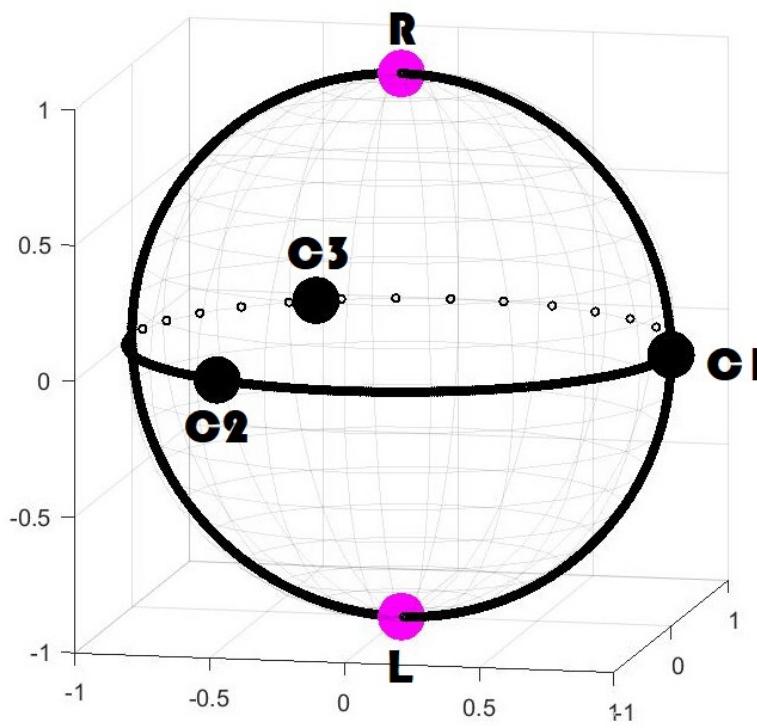


$$\bar{\Gamma} = \Gamma_{\text{bare}} - \Gamma_{\text{ref}}$$

to appear,
w. Dandekar,
Mazumdar, Lederer

Novel derivation of Lagrange's equilateral soln & its stability

Equilibria in \vec{g} \equiv rigid rotation



\vec{g} space

Perspective

Newton 1687

$$\ddot{\vec{r}}_1 = -G \left(\frac{m_2}{r_{12}} + \frac{m_3}{r_{13}} \right)$$

Lagrange 1772

$$\frac{d^2}{dt^2} r_{12}^2 = \dots$$

...

Jacobi 1842

Elimin. of nodes

Levi-Civita 1915

$$\vec{\omega}, \vec{H}$$

Murnaghan 1936,

planar symmetric r_{ab}^2

(3 d.o.f)

Van Kampen & Winther 1937

3d

(4 d.o.f)

Lemaître 1952 symmetric & smooth at collinear

Moeckel & Montgomery 2013

planar, shape sphere,
modernize

Helium spectrum

$$m_1 = m_2 = m_e \quad M = \infty$$

Hylleraas 1929 $\vec{J}^2 = 0$ \sim d.o.f : r_{ab}^{-2}

Gronwall 1932 \sim $\vec{g} \in \mathbb{R}^3$

Fock 1954 \sim $w \in \mathbb{R}^4$

Tkackenko 1978 , H for general \vec{J} .

Novelties here

- * Defn. of w
- * Formulation in \vec{J} .

Recap.

